

Investigation of Synthetic Aperture Methods in Ultrasound Surface Imaging Using Elementary Surface Types

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Abstract

Synthetic aperture imaging methods have been employed widely in recent research in non-destructive testing (NDT), but uptake has been more limited in medical ultrasound imaging. Typically offering superior focussing power over more traditional phased array methods, these techniques have been employed in NDT applications to locate and characterise small defects within large samples, but have rarely been used to image surfaces. A desire to ultimately employ ultrasonic surface imaging for bone surface geometry measurement prior to surgical intervention motivates this research, and results are presented for initial laboratory trials of a surface reconstruction technique based on global thresholding of ultrasonic 3D point cloud data. In this study, representative geometry artefacts were imaged in the laboratory using two synthetic aperture techniques; the Total Focusing Method (TFM) and the Synthetic Aperture Focusing Technique (SAFT) employing full and narrow synthetic apertures, respectively.

Three high precision metallic samples of known geometries (cuboid, sphere and cylinder) which featured a range of elementary surface primitives were imaged using a 5MHz, 128 element 1D phased array employing both SAFT and TFM approaches. The array was manipulated around the samples using a precision robotic positioning system, allowing for repeatable ultrasound derived 3D surface point clouds to be created. A global thresholding technique was then developed that allowed the extraction of the surface profiles, and these were compared with the known geometry samples to provide a quantitative measure of error of 3D surface reconstruction. The mean errors achieved with optimised SAFT imaging for the

29 cuboidal, spherical and cylindrical samples were 1.3 mm, 2.9 mm and 2.0 mm respectively,
30 while those for TFM imaging were 3.7 mm, 3.0 mm and 3.1 mm, respectively. These results
31 were contrary to expectations given the higher information content associated with the TFM
32 images. However, it was established that the reduced error associated with the SAFT
33 technique was associated with significant reductions in side lobe levels of approximately
34 24dB in comparison to TFM imaging, although this came at the expense of reduced
35 resolution and coverage.

36

37 *Keywords:* Ultrasound imaging, Total Focussing Method, Synthetic Aperture Focussing
38 Method, Full Matrix Capture, Robotics.

39

40 **1. Introduction**

41 A number of robotically guided knee arthroplasty systems require a preoperative 3D model of
42 the joint, with which the surgery can be planned and implemented [1]–[4]. Computed
43 Tomography (CT) is seen as the ‘gold standard’ in this area, but it is costly [5] and can apply
44 a dose of ionising radiation greater than the yearly background dose of 2.2 mSv [6], which is
45 potentially dangerous to the patient [7]. Ultrasound imaging has the potential to provide an
46 alternative to CT in this capacity by offering comparable accuracies, while reducing cost and
47 eliminating the risk associated with ionising radiation.

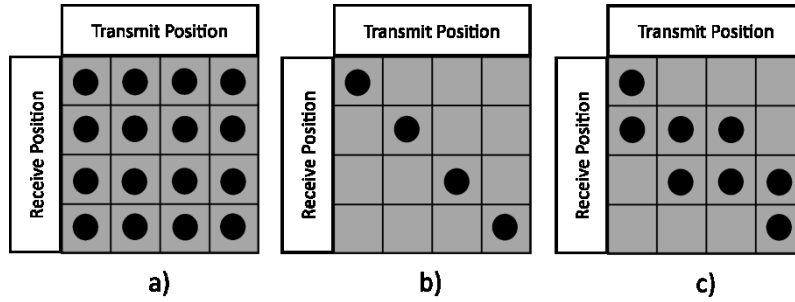
48 Synthetic aperture imaging methods have become commonplace in research in non-
49 destructive testing (NDT), allowing for improved focussing capability and increased
50 resolution over more traditional B-scan methods [8]. With these attributes, such techniques
51 could provide an improvement over traditional medical imaging methods in accurately
52 reconstructing the bony surfaces of the knee. However, these techniques have found little
53 uptake in medical ultrasound imaging, with standard commercial systems lacking the
54 versatility to perform Full Matrix Capture (FMC) [9] – a requirement of popular
55 reconstruction algorithms such as the Total Focussing Method (TFM). Additionally, real
56 time, high resolution imaging is usually a requirement of medical ultrasound systems. This is
57 difficult to achieve using synthetic aperture methods, in that the techniques are inherently
58 computationally expensive. As such, high frame rates often cannot be achieved, even when
59 exploiting the parallelisable nature of the calculations [10].

60 Research into synthetic aperture methods in NDT has, for the most part, concentrated on
61 locating and characterising small defects within relatively large samples [11]–[13]. While
62 TFM has been used for surface imaging in dual media compensation calculations [14], [15],
63 these and other auto focussing techniques have been limited to relatively simple, continuous
64 surfaces [16], [17]. Imaging complex surfaces such as bones would, on the other hand,
65 require surface reconstruction of highly variable and often discontinuous surface types [18].
66 A further challenge is to fully represent the entire surface under inspection, requiring, firstly,
67 a high number of images and, secondly, an accurate probe positioning system.

68 To meet these challenges and to test the ability of synthetic aperture methods to reconstruct
69 surfaces at a fundamental level, three high precision metallic samples of known, simple
70 geometries were imaged. These samples provided elementary surface types and features that
71 would be found in complex samples, such as bone. Data was captured using FMC and was
72 processed using both TFM and a form of the Synthetic Aperture Technique (SAFT). These
73 methods provided a contrast between the two extremes of the spectrum in synthetic aperture
74 methods – TFM comprising a fully populated transmit-receive matrix and SAFT a minimally
75 diagonally populated transmit-receive matrix. As can be seen in Fig. 1, TFM offers the
76 maximum possible focussing power by synthetically focussing using the full aperture in
77 reception. SAFT methods, on the other hand, employ a sub aperture which diminishes the
78 focussing capabilities and imposes a higher level of positional dependence on the
79 reconstruction. High data throughput was achieved using an FPGA-based phased array
80 controller, while high speed data processing times were made possible through a Graphics
81 Processing Unit (GPU) implementation of the synthetic aperture algorithms [19]. The probe
82 was manipulated using a robotic precision positioning system, which provided accurate
83 positional data, allowing for ultrasound derived 3D surface point cloud reconstruction. The
84 3D reconstructed surfaces were compared to the known reference models and the
85 performances of the imaging methods were compared.

86

87



88

89 Fig. 1: Transmit-receive matrices of a four element array for TFM (a), single element receive
 90 SAFT (b) and multiple element receive SAFT (c).

91

92 The structure of the paper is, firstly, an introduction to the data capture methods and imaging
 93 algorithms employed, which is then followed by a description of the experimental apparatus
 94 and method. The results for each sample are then presented and discussed separately, after
 95 which a discussion of performance limitations is presented. This is followed, finally, by a
 96 summary and conclusions.

97

98 2. Materials and Methods

99 2.1 Synthetic Aperture Methods

100 The concept of FMC [8] is to excite one element of the phased array and receive on all the
 101 others. The succeeding element is then fired and all elements become receivers once again.
 102 This process is repeated for all N elements, producing a $N \times N$ matrix of time signals, which
 103 is known as the full matrix.

104 While real time synthetic aperture implementations have been achieved [10], [20], high
 105 resolution images created using probes with high element counts must still be produced in
 106 post processing. With the full matrix, it is possible to apply numerous processing methods to
 107 the same data set. One popular method is TFM, which employs every element in the array (ie.
 108 the full aperture) to synthetically focus in transmission and reception for every pixel in the
 109 image. This process begins by discretising the image region into a grid of points, each of
 110 which defines the location of a pixel in a scalar image. The intensity of a particular pixel can
 111 be calculated using Equation 1, in which $S_{i,j}$ is the time-trace associated with a transmission

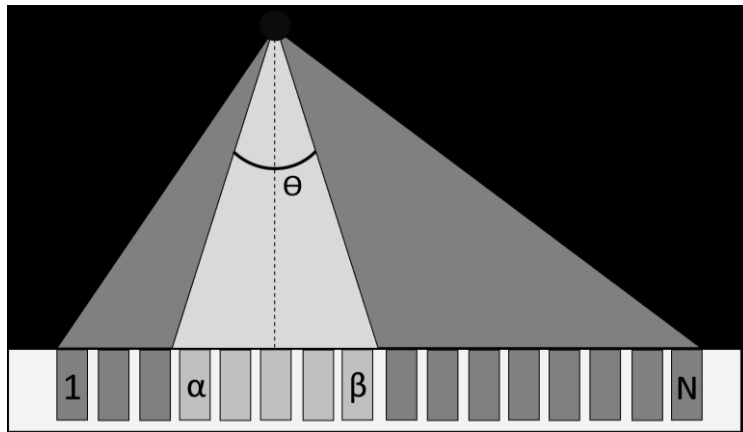
112 at the i th element and reception at the j th element, while (x, y) are the coordinates of the
 113 pixel. The time of travel from the transmitting element to the pixel is represented by $T_{i(x,y)}$,
 114 while that from the pixel to the receiving element is signified by $T_{j(x,y)}$. This is summed over
 115 the number of elements in the array, N , and is then repeated for every pixel in the image.

$$I(x, y) = \sum_{i,j=1}^N S_{i,j}(T_{i(x,y)} + T_{j(x,y)}) \quad (1)$$

116

117 In addition to TFM, a form of SAFT was employed in which the same process was carried
 118 out, but focussing was not performed with the full aperture. Instead, the elements constituting
 119 the synthetic aperture were defined by the position of the pixel in question. As shown in Fig.
 120 2, while the TFM aperture included every element, the SAFT aperture was restricted to those
 121 elements contained within an isosceles triangle defined by the angle Θ .

122



123

124 Fig. 2: Graphical representation of the TFM and SAFT synthetic aperture definitions, with
 125 the elements constituting the SAFT aperture shown in a lighter shade.

126

127 The intensity of a pixel in the image is, then, given by:

$$I(x, y) = \sum_{i=1}^N \sum_{j=\alpha}^{\beta} S_{i,j}(T_{i(x,y)} + T_{j(x,y)}) \quad (2)$$

128

129 where α and β are the first and last elements of the aperture, as displayed in Fig. 2. These
130 values vary depending on the definition of Θ , with an increase in the angle increasing the
131 aperture size.

132 Different SAFT aperture widths were trialled at 10° intervals from 10° to 180°, allowing for
133 an assessment of performance and characteristics that was representative of the full spectrum
134 of possible aperture widths. While different surface types affected performance, the 20°
135 aperture SAFT was found to offer the greatest contrast in performance to the full aperture
136 reconstruction of TFM whilst maintaining the ability to reconstruct most surfaces.

137

138 *2.2 Experimental Apparatus*

139 Three representative surface geometry forms were considered to encompass interaction with
140 curves, flat surfaces and edges which are the primitive geometry forms encountered in bone
141 geometries associated with the knee. In order to test these surface types both in isolation and
142 in combination, three samples were prepared, the first of which was a brass sphere with a
143 diameter of 25.0 mm (Dejay Distribution Ltd., Cornwall, UK). In addition to this, a cuboidal
144 sample was manufactured in aluminium, with dimensions 25.1 mm × 35.4 mm × 35.2 mm.
145 Finally, a cylindrical aluminium sample was produced with a height of 62.5 mm and a
146 diameter of 50.0 mm. The cylindrical sample was, additionally, flattened off 15.0 mm from
147 the centre. In doing this, the sample provided all three surface features, as can be seen in Fig
148 3. To ensure the majority of the surface was accessible during inspection, a mount was
149 manufactured that allowed the samples to be elevated, such that the interrogation array probe
150 had good line of sight access to the whole of the samples.

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152

153

Fig. 3: The spherical, cuboidal and cylindrical samples shown with the mount.

154

155 A 128 element Vernon (Tours, France) phased array with a centre frequency of 5 MHz was
156 employed for ultrasonic acquisition. This offered both a large aperture width of 89.6 mm and
157 a theoretical resolving limit of less than 0.2 mm in water. Standard preoperative CT scans
158 employ 1.0 mm thickness slices [1], [21]–[23], limiting the resolution to 1.0 mm. The probe
159 was, therefore, significantly within the required spatial resolution limit, relative to reference
160 CT imaging.

161 The array was excited and interrogated by a FlawInspecta (Diagnostic Sonar Ltd., Livingston,
162 UK) phased array controller. The platform is modular, with parallel digitisation achieved
163 using FlexRIO FPGA cards. The configuration used herein employed two 32 channel
164 digitisers, allowing for parallel reception of 64 elements, with a maximum of 4 elements in
165 simultaneous transmission. Therefore, two firings were required for each transmission event
166 when using all 128 elements. This was the only limiting factor on the Pulse Repetition
167 Frequency (PRF) originating from the hardware, with the only other constraint being that of
168 wave travel within the material [24]. However, at the time of data capture, the firmware was
169 not optimised, meaning that a bottleneck was created in data transfer [25] which resulted in a
170 variable frame rate of approximately 0.3 Hz. Using a the speed of sound in water (1480 ms^{-1}
171 [26]), images with a width of 13.42 cm and a depth of 8.95 cm were reconstructed.

172 In order to reconstruct 3D surfaces using 2D images, accurate probe manipulation and
173 positional recording were vital. To this end, a KUKA KR 5 arc HW industrial robot was

174 employed, providing six degrees of freedom and the ability to implement a range of array
175 probe paths. Industrial robotic manipulators are seeing increasing application in high
176 precision manipulation tasks [27]–[29] despite known issues with absolute accuracy. Despite
177 this shortcoming, such robots are extremely repeatable in position, and have great advantages
178 in being able to move in complex curved paths with 6 degrees of freedom whilst maintaining
179 constant standoff and normality to the local surface geometry. Additionally the ability to
180 employ CAD/CAM based off line programming allows ease of programming to produce
181 complex tool paths specific to a part with known geometry [30], [31].

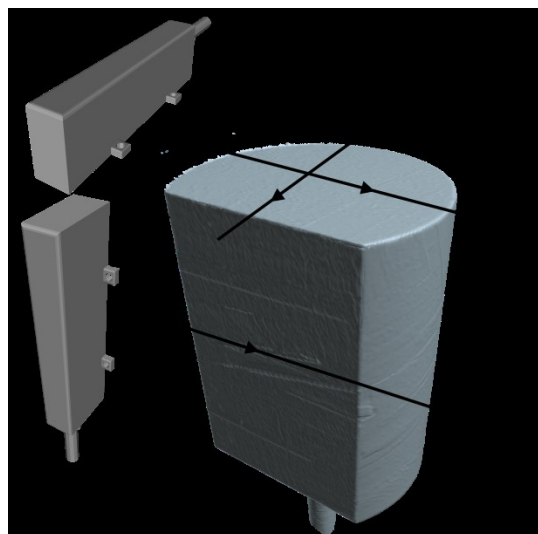
182 The correct calibration of the inspection probe tool centre point (TCP) was critical to the
183 attained accuracy, as small errors in this physical position with respect to the flange,
184 translated into much larger errors in absolute position as the probe was located at the end of
185 the physical kinematic chain. A standard two stage KUKA TCP calibration method was
186 employed [32]. The first phase of the procedure defined the position of the TCP relative to
187 the flange and involved moving the origin of the tool reference frame to a static reference
188 point four times, each with a different robot pose. This was achieved using a spike which was
189 manufactured such that the position of the tip relative to the flange corresponded to the centre
190 of the probe face. The second stage allowed for tool orientation calibration, which involved
191 moving points on the X axis and the X-Y plane of the tool coordinate system to the reference
192 point. Two further calibration spikes were employed to accomplish this, which conformed to
193 the described requirements. The X-Y plane was defined such that it corresponded to the 2D
194 imaging plane associated with the probe. Two bespoke probe holders were manufactured; one
195 of which was parallel to the 6th axis of the robot, while the other was perpendicular, allowing
196 for full line of sight access to the sample while maintaining full probe submersion. Each
197 probe mount was calibrated as described above, producing calibration errors of 0.6 mm and
198 0.7 mm, each below the recommended industry maximum error of 0.8 mm.

199 To independently assess the true position of the probe compared to the reported KUKA
200 position during manipulation, six Vicon T160 cameras (Vicon Motion Systems, Oxford, UK),
201 were employed. A marker set of five 12 mm retro reflective markers allowed for the TCP
202 position to be tracked while the TCP followed a hemispherical scan path with a similar
203 working volume to that required to image the samples. The path required that a range of
204 complex poses be adopted by the robot, which imposed highly varied joint angle
205 combinations. The KUKA TCP position was relayed to the host PC every 12 ms via ethernet,
206 using the Robot Sensor Interface (RSI) software [33]. The corresponding, temporally

207 synchronised Vicon measurement was attached to each KUKA measurement. The resulting
208 point clouds were matched using Iterative Closest Point (ICP) to compensate for the
209 difference in the origin and orientation of the coordinate systems of the two systems. The
210 absolute Euclidean distances between each of the corresponding KUKA and Vicon derived
211 positions were then calculated, resulting in a mean error of 0.5 mm. This was considered a
212 worst-case error, as the path was highly complex which would be expected to produce larger
213 errors than would be found in more simple paths. While optical tracking is known to produce
214 relatively large errors, these are dependent on the size of the measurement volume [34]. As
215 the measurement and calibration volumes were small (0.002 m^2 and $\sim 6 \text{ m}^2$), these errors were
216 minimal.

217 In the current application, there is no a-priori information available that would allow for
218 probe trajectory path planning for the geometry of the sample to be scanned. Indeed in the
219 ultimate desired application in knee joint imaging, all that would be available would be the
220 rough working volume definition around the subject's limb. Therefore, complex path
221 programming was not required or indeed possible, and a simple rectilinear scan path around
222 each object was employed. This had the advantage of simplification of programming using
223 the inbuilt KUKA Robot Language (KRL) – a BASIC-like, domain specific language which
224 allowed for simple tool paths to be defined. Such a path is shown in Fig. 4 for inspection of
225 the cylindrical sample. The path, while not maintaining normality or a specific standoff,
226 provided full coverage of the surface.

227



228

229

Fig. 4: Cuboidal tool path written in KRL.

230

231 The samples were placed in a water bath to allow for complete submersion, as can be seen in
232 Fig. 5. The KRL code describing the path required initial coordinates to be defined. These
233 were identified by manually positioning the probe to a point at which the probe face would
234 remain submersed at all times. Additionally, the length, width and height of the cuboid
235 defining the path were altered with each sample, so that a minimum standoff of at least 20
236 mm was maintained.

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238

239 Fig. 5: Experimental setup, showing the submerged sample, with the probe face fully
240 submerged.

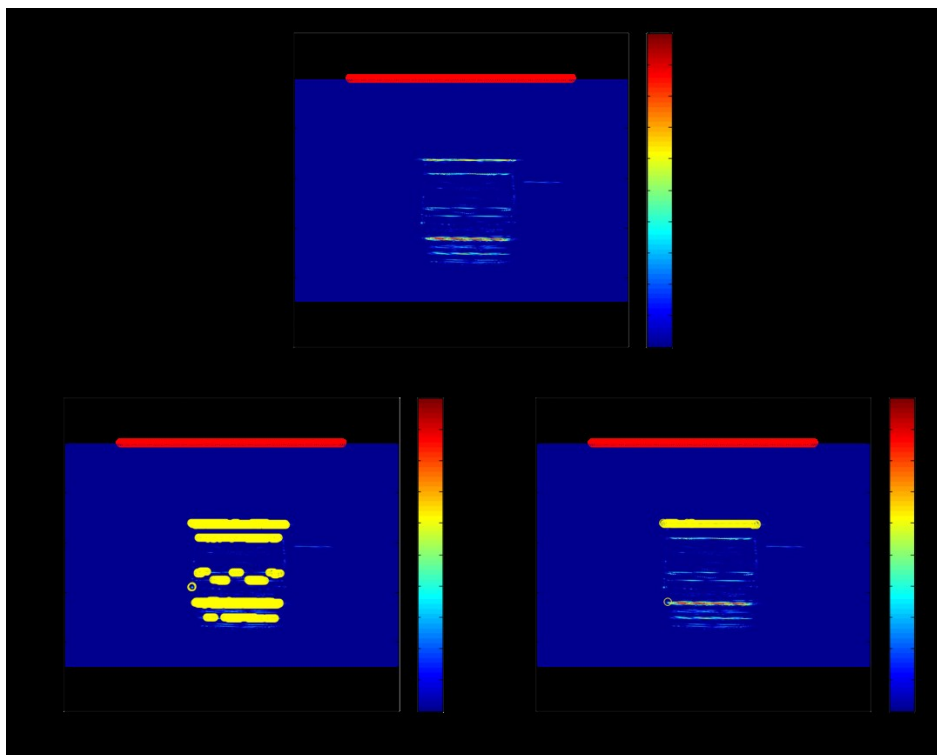
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242 *2.3 Image and Surface Point Cloud Construction*

243 The process of constructing images using synthetic aperture methods is computationally
244 expensive, but is highly parallelisable [20], [35]. A software platform, cueART, has been
245 developed by the Centre for Ultrasonic Engineering (CUE) at the University of Strathclyde,
246 which allows for significant reductions in computation time by implementing the algorithms
247 on a GPU [14], [19]. In doing so, cueART allowed for high resolution images to be produced
248 using 128 element FMCs in a practical time frame. An image depth of 8

249 To identify the surface profile in each image, global thresholding was employed, which
250 revealed the coordinates of all the pixels with intensities above the defined decibel limit.
251 However, the images did not display only the first surface reflection, but multiple others, as
252 can be seen in Fig. 6 (a). The first surface was the region of high intensity closest to the probe
253 face (signified by circles), which represented the true location of the outer surface of the
254 sample. The second surface, directly beneath the first, was caused by reflections from the
255 back wall of the sample. The most prominent of the surface profiles in the image, seen at the
256 bottom of the image, was caused by the reflected waves from the first surface reflecting on
257 the probe face and making the return journey. As such, the line was twice the distance from
258 the probe face as the true surface representation. In addition to these, there were numerous
259 other false interface indications caused by further back wall reflections.

260



261

262 Fig. 6: A typical SAFT image obtained from the inspection of the cuboidal sample (a).
263 Thresholding alone extracted all the erroneous surface representations present in the image
264 (b), while the surface extraction algorithm isolated the true surface (c).

265

266 When global thresholding was performed on such an image, all the areas of high intensity
 267 were identified, as can be seen in Fig. 6 (b). In order to isolate the coordinates of the first,
 268 true surface, all the coordinates above the threshold were stored in a matrix. Any coordinates
 269 which contained the same X value were discarded, with the exception of that with the
 270 smallest Z value. As such, it was ensured that for every column of pixels, only the pixel
 271 closest to the probe face and above the decibel limit would be recorded. The result of this can
 272 be seen in Fig. 6 (c), where the erroneous surfaces have been eliminated.

273 The 2D coordinates representing surface contours from each image were recorded relative to
 274 the centre of the probe face. To place the points in 3D space relative to the coordinate system
 275 of the robot, the coordinates were rotated then translated using the corresponding measured
 276 KUKA position and orientation. This is presented in equation (3), where the subscripts f , o
 277 and T represent the final 3D coordinates, the original coordinates and the TCP coordinates
 278 needed for translation, respectively. Additionally, R_{xyz} is the rotation matrix in x, y, z order.
 279 Coordinate conversion from 2D to 3D was simple, as the TCP was calibrated such that the
 280 position and orientation corresponded directly to the imaging plane.

$$\begin{pmatrix} X_f \\ Y_f \\ Z_f \end{pmatrix} = R_{xyz} \begin{pmatrix} X_o \\ Y_o \\ Z_o \end{pmatrix} + \begin{pmatrix} X_T \\ Y_T \\ Z_T \end{pmatrix} \quad (3)$$

281

282 Finally, the surface point clouds were imported into CloudCompare (EDF/Telecom
 283 ParisTech, Paris, France) – an open source software package designed for comparing 3D
 284 point cloud data. The point clouds were matched with the reference models, firstly, by
 285 manually manoeuvring the cloud so that an approximate match was achieved. The point
 286 cloud was then finely matched by way of ICP. This was necessary step as the position and
 287 orientation of the samples was unknown relative to the KUKA coordinate system. While it
 288 would be advantageous to accurately position the sample such that the true position of the
 289 sample would be known relative to every packet of KUKA positional data, the samples were
 290 positioned manually and approximately to more realistically produce the final application,
 291 where no such reference data would be available from which to gain absolute positional data.
 292 It should also be said that while CloudCompare provided a convenient platform on which to
 293 compare point cloud data and visualise results, it was not vital to the system, as it did not
 294 contribute to the reconstruction process and its processes could be simply replicated on other
 295 platforms.

296 The error for each point in the ultrasound-derived point cloud was then calculated by finding
297 the absolute Euclidean distance between the point and the nearest vertex on the surface of the
298 reference model after matching. From these values, mean error, maximum error and standard
299 deviation were calculated for each ultrasound-derived surface point cloud.

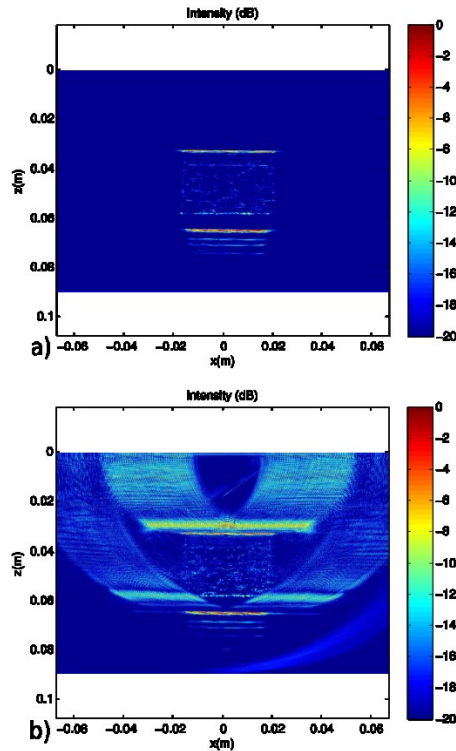
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301 **3. Results and Discussion**

302 *3.1 Cuboidal Sample*

303 A pronounced difference can be seen in Fig. 7 between the images produced by SAFT and
304 TFM. The SAFT-derived image shows a clear outer surface representation, along with other
305 spurious reflections, as described in section 2.3. While the TFM image displays the same
306 surface representations, it also exhibits side lobes of high intensity. These side lobes were
307 artefacts generated in the reconstruction algorithm and attributed to true reflectors [19]. The
308 typical maximum side lobe intensity for the side lobes found in the TFM images was -6 dB,
309 while that for the SAFT images was -30 dB. The consequence of this was that when contour
310 extraction was employed, the true surface was often not identified for both SAFT and TFM
311 reconstructions, with the artefacts above it instead being extracted. In an effort to limit the
312 effect of this, the threshold limit of the TFM images was set to -5 dB, while that for the SAFT
313 images was -12 dB (the difference reflecting the side lobe levels). In doing this, the number
314 of pixels representing the surfaces was reduces, therefore resulting in an undesirable
315 reduction in coverage.

316 Additionally, it can be seen in Fig. 7 (a) that the lower surface representation appears to
317 display a higher intensity than the upper, true surface. The reason for this is that the pixels
318 with a larger Z coordinate would have employed a greater number of receiving elements
319 during reconstruction, implying a greater scalar intensity. As can be seen in Fig. 2, as the
320 pixel moves further from the probe face, the angle, Θ , remains constant, the width of the base
321 of the triangle increases and, therefore, the number of elements in the receiving aperture
322 increases.



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Fig. 7: SAFT (a) and TFM (b) images resulting from a typical FMC from the cuboidal sample.

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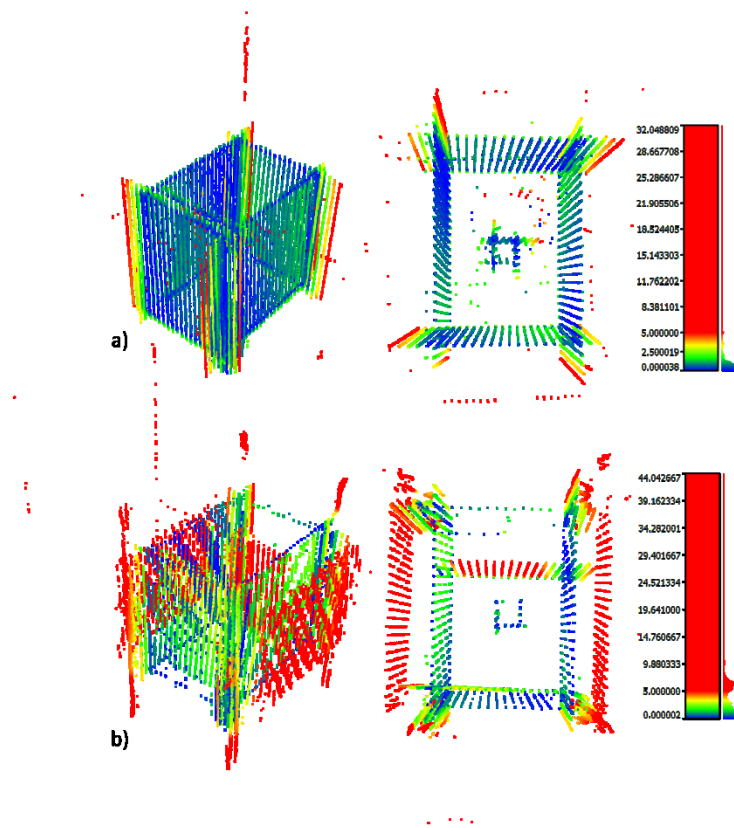
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Once constructed and compared, the SAFT point cloud achieved significantly lower errors than TFM, as can be seen in Table 1. Additionally, the TFM point cloud was not as dense as the SAFT counterpart, which was expected, given the higher threshold limit. Further, the TFM cloud appeared to contain multiple surfaces, as can be seen in Fig. 8 (b). The first cause of this was high intensity side lobes, while the second cause was extraction of incorrect surfaces in each image, which was made possible by the higher threshold limit associated with TFM. Comparing parts (a) and (b) of Fig. 8, it can be seen that SAFT processing provided a significantly more accurate depiction of the sample. As can be seen in Fig. 8 (a), a cuboid-like structure is present, with most of the points in these regions achieving sub-millimetre accuracy. However, there were outlying points which were a result of noise, registered due to the low threshold limit used in the SAFT images. In addition to this, each of the faces extended beyond the edge, which increased the mean error and standard deviation significantly.



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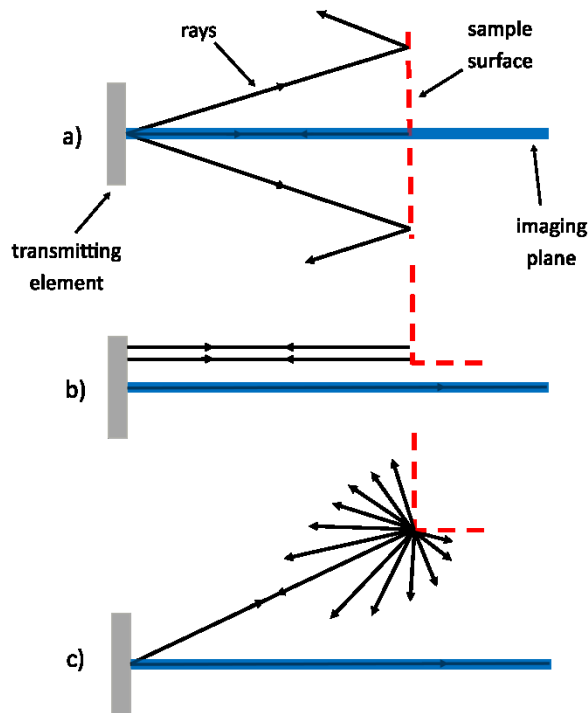
341 Fig. 8: SAFT (a) and TFM (b) derived surface point cloud reconstructions of the cuboidal
 342 sample, showing errors at two viewing angles.

343

344 The cause of this was not positional inaccuracy on the part of the robot controller or in the
 345 point cloud construction procedure, but rather the physical nature of the ultrasound beam
 346 itself. The beam shape of an individual element of the array has an associated thickness, not
 347 only in the width along the direction of the full aperture, but also associated with the
 348 elevation of each element. This beam width is not considered during 3D construction, as the
 349 images are regarded as 2D.

350 The effect of this beam width is usually negligible when imaging planar surfaces at right
 351 angles. The reason for this is that for relatively polished surfaces, such as that employed
 352 herein, specular reflection dominates. Therefore, any off-axis ultrasonic energy is directed
 353 away from the receiving element and only the energy within the imaging plane is recorded, as
 354 illustrated in Fig. 9 (a).

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Fig. 9: Different reflection types from a planar surface (a, b) and an edge (c).

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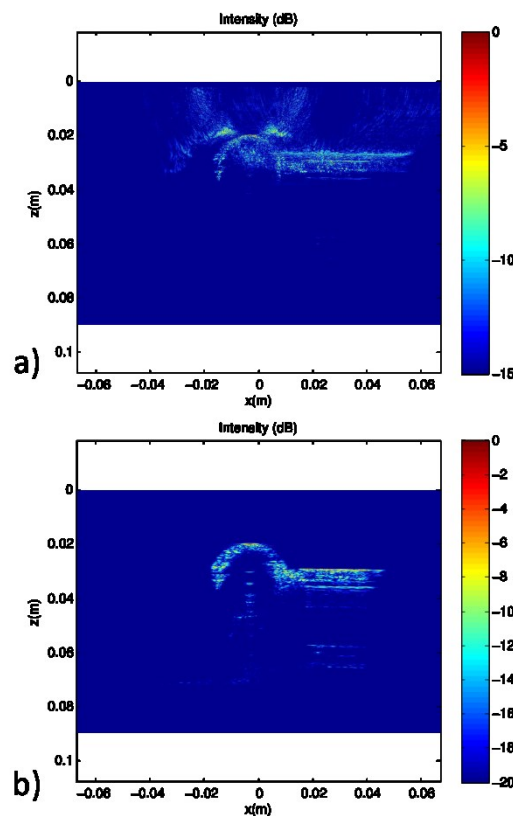
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3.2 Spherical Sample

However, when the imaging plane initially passes the edge of a sample, the main beam from the element still reflects from the planar surface. As such, the resultant image will display a surface, but this will be inaccurately placed due to the assumption of a 2D imaging plane positioned through the centre of the element, as demonstrated in part (b) of Fig. 9. Given that the elevation of the elements is 10 mm, this effect would account for surfaces being registered at most 5 mm past the edge of the sample. However, in the SAFT-derived point cloud (the most accurate of the two), surfaces were reconstructed over 9 mm from the edge of the sample. This was caused by reflections from the edge of the sample, resulting from off-axis transmissions. While specular reflection dominates in planar surfaces, in the case of an edge, diffuse reflection dominates, as illustrated in Fig. 9 (c). While the intensity of the signal returned in this instance would be lower than that from the planar surfaces, the images were evaluated in a decibel scale on an individual basis, meaning the overall lower amplitude of the signals would have little impact.

374 The TFM images resulting from the scan of the spherical sample displayed significant side
375 lobes, as seen in Fig. 10 (a). The most prominent of these formed as two islands of intensity
376 similar to those that would be expected in from a point reflector. As can be seen in Fig. 10
377 (b), while the SAFT image showed no side lobes, there was a reduction in the resolution of the
378 surface profile relative to the TFM image, with a thickening of the surface representation.
379 This was a result of a reduction in the resolving power associated with SAFT relative to
380 TFM, which offered the maximum possible resolving power [36]. Additionally, TFM allowed
381 for reconstruction of more of the surface than SAFT. This was because the narrow synthetic
382 aperture of SAFT only allowed for reconstruction of surfaces directly below the probe face.

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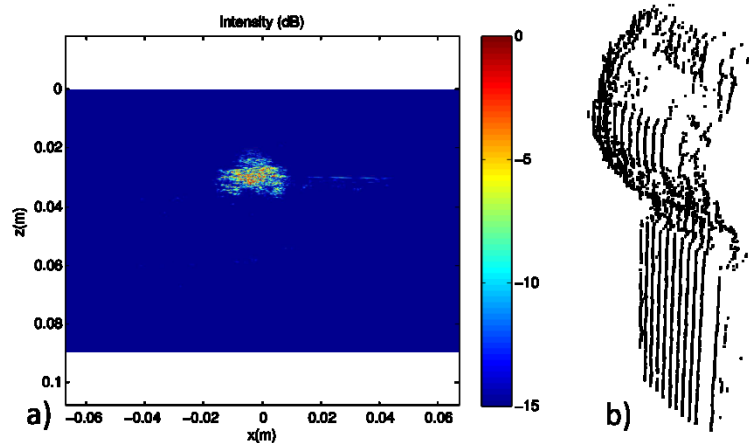
385 Fig. 10: TFM (a) and SAFT (b) images resulting from a typical FMC from the spherical
386 sample.

387

388 When the probe moved away from the centre of the sphere, the obvious surface profiles
389 exemplified in Fig. 12 were not present. Instead, the profiles became distorted, as can be seen
390 in Fig. 11 (a). The surface extraction and 3D reconstruction of these images resulted in a lack

391 of curvature in both the sphere and the supporting rod. This is demonstrated in Fig. 11 (b),
392 which shows the point cloud resulting from one straight scan line of the full scan. It would be
393 expected that the points would display an obvious curvature, but they instead possessed an
394 almost complete lack of curvature. The same effect was encountered with the mount in the
395 cuboidal sample, as seen in Fig. 8.

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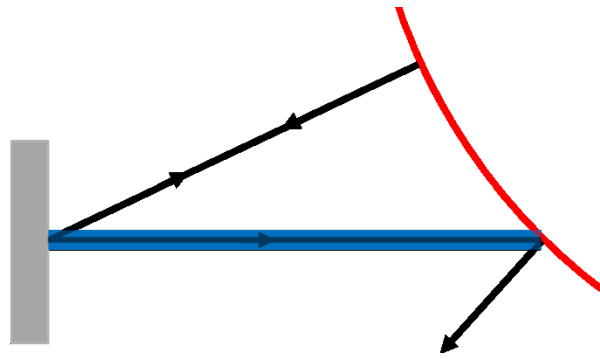
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398 Fig. 11: SAFT image from FMC captured off the central axis of the sphere (a) and the surface
399 reconstruction of one of the scan lines (b).

400

401 The distortion effect was caused by the three dimensional nature of the transmitting beam. As
402 illustrated in Fig. 12, the rays radiating from the image plane were reflected away from the
403 receiver as the probe moved away from the centre of the sphere. However, the rays outside
404 the image plane were received and placed inaccurately. Numerous ray paths along the face of
405 the transducer caused many reflections to be received at different times, resulting in the
406 distortion effect. As the surface extraction algorithm discriminated in favour of those pixels
407 closest to the probe face, the reflections originating close to the centre of the sphere were
408 always favoured.

409



410

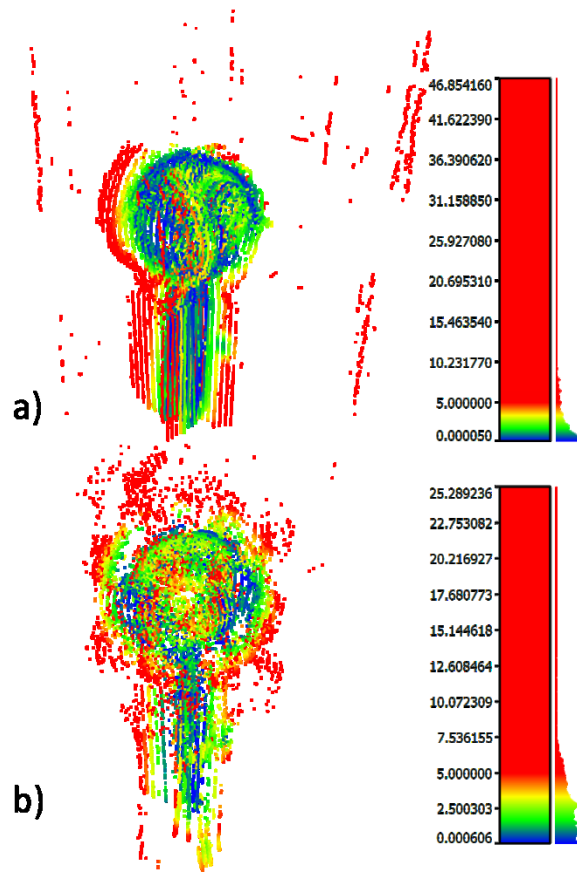
411

Fig. 12: Specular reflection resulting from a curved surface.

412

413 These effects had a significant impact on the reconstructed surface point clouds, with an
414 increase in the SAFT errors relative to the cuboidal results, as shown in Table 1. While the
415 SAFT point cloud maintained a lower mean error than the TFM cloud, it produced a
416 significantly higher standard deviation and maximum error. It can be seen in Fig. 13 that the
417 effects of specular reflection from the curved surface were more pronounced in the SAFT
418 data than the TFM, causing the higher levels of error. The TFM point cloud also had a large
419 number of inaccurate points, but they were of a different nature, being relatively close
420 proximity to the true surface. This was, once again, caused by significant side lobes in the
421 TFM images.

422



423

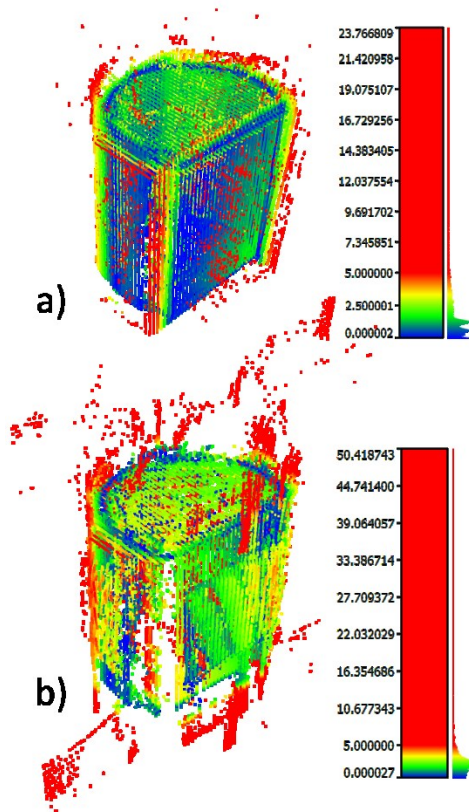
424 Fig. 13: SAFT (a) and TFM (b) derived surface point cloud reconstructions of the spherical
 425 sample.

426

427 3.3 Cylindrical Sample

428 The cylindrical sample included all three representative surface primitives (curves, flat
 429 surfaces and edges), providing the opportunity to assess them in combination. The effect of
 430 edges can be seen in Fig. 14 (a), where, towards the left of the image, inaccurate lines of
 431 points can be seen extending past the edges of the sample. The flat surfaces were generally
 432 accurate. However, it can be seen in part (a) that most of the flat surface at the top of the
 433 sample had errors of approximately 1 mm. This was most likely caused by the matching
 434 procedure, which minimises the error for all points, rather than merely those of high
 435 accuracy.

436



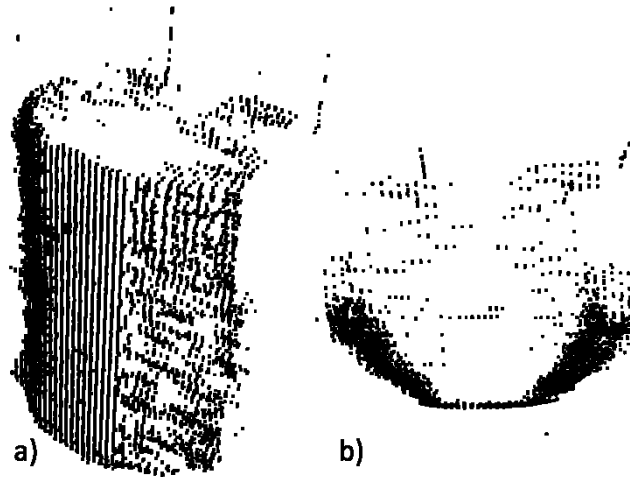
437

438 Fig. 14: SAFT (a) and TFM (b) derived surface point cloud reconstructions of the cylindrical
 439 sample.

440

441 The effect of the curved surfaces is detailed in Fig. 15, where, around the centre of the curve,
 442 there appears to be a complete lack of curvature. This is similar to the features seen in Fig. 11
 443 (b), the cause of which was the specular effects explained in section 3.2. In addition to this,
 444 when the probe moved farther from the centre of the curve, it can be seen that the lines
 445 became erratic and sparse. This was due to the distorting effect described in section 3.2,
 446 which did not provide obvious surface profiles to be extracted. The effects of curvature were
 447 prominent due to the positioning of the probe relative to the surface. If a path was employed
 448 which maintained normality to the sample surface, the effects would be reduced significantly.

449



450

451 Fig. 15: Surface reconstruction of the scan line that moves past the side opposing the flat
 452 surface of the cylindrical sample.

453

454 Table 1: Results of comparison with the reference models.

Sample	Processing Method	Mean Error (mm)	Maximum Error (mm)	Standard Deviation (mm)
Cuboid	TFM	3.7	44.0	3.1
Cuboid	SAFT	1.3	32.0	1.8
Sphere	TFM	3.0	25.2	2.6
Sphere	SAFT	2.9	42.9	4.4
Cylinder	TFM	3.1	52.2	4.5
Cylinder	SAFT	2.0	26.5	2.1

455

456

457 *3.4 Discussion*

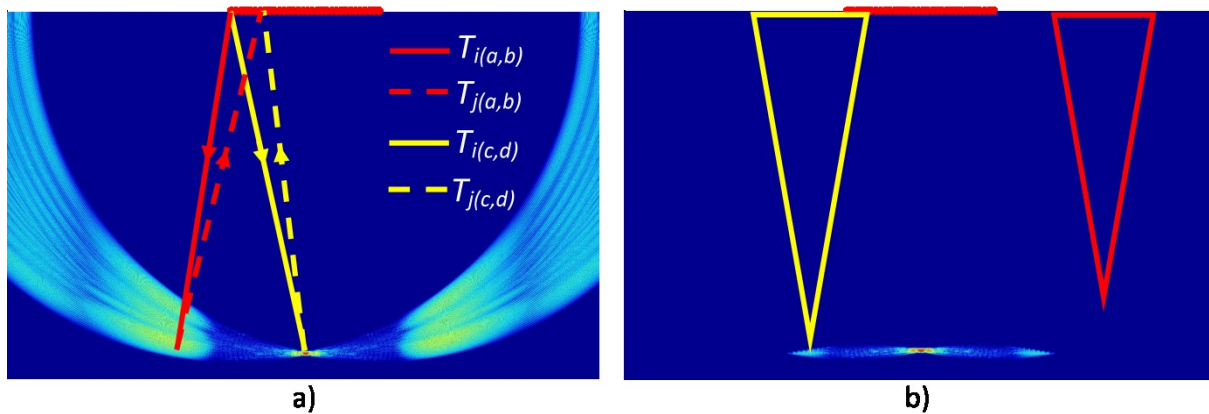
458 The images constructed using TFM have had noticeably higher side lobe content, leading to
 459 an increase in the associated measured surface error for all three sample geometries
 460 investigated. In traditional ultrasound imaging, the term “side lobes” refers to imaging
 461 artefacts that result from regions of ultrasonic energy which are produced off-axis relative to
 462 the main lobe during transmission [37]. In synthetic aperture methods, however, side lobes

463 are regions of high intensity not attached to the main lobe [19] and are a result of the image
 464 construction algorithm itself.

465 This can be explained by considering a point spread function generated using TFM.
 466 Employing a ray-based model based on that described in [8], FMC data was simulated,
 467 providing the response of a point reflector 10 mm from the face of the probe. The medium
 468 was defined with a longitudinal speed of sound of 1480 ms^{-1} , the array with 32 elements with
 469 a pitch of 0.7 mm and the sampling frequency as 100 MHz. The output of each element was
 470 modelled as a 5 cycle, Gaussian windowed tone burst with a centre frequency of 5 MHz and a
 471 -6 dB bandwidth of 50%, as has been typically employed before [8], [38], [39].

472 As can be seen in Fig. 16 (a), as well as the obvious point reflector representation at (a,b),
 473 there are significant side lobes located either side. The reason these occur is because the
 474 algorithm does not discriminate based on pixel location. For example, for transmitting
 475 element i and receiving element j , the time of flight to pixel locations (c,d) – the position of
 476 the point reflector – and (a, b) is the same. In other words, with reference to (1), $S_{i,j}(T_{i(a,b)} +$
 477 $T_{j(a,b)}) = S_{i,j}(T_{i(c,d)} + T_{j(c,d)})$. As such, the contribution to the intensity of the pixel, for that
 478 particular transmit-receive pair will be the same at both locations.

479

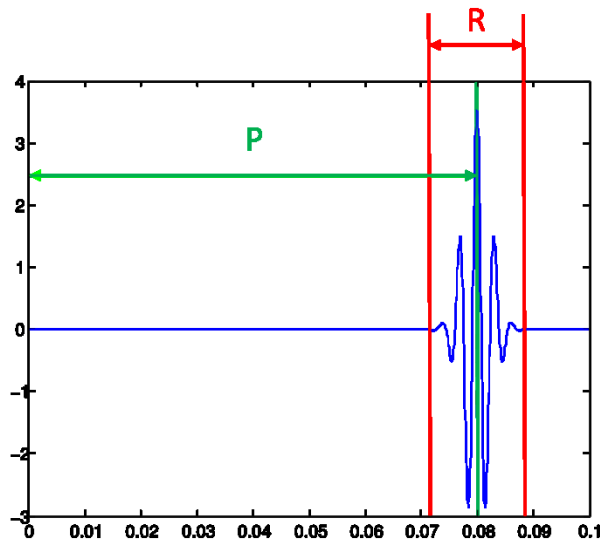


480

481 Fig. 16: Part (a) shows a TFM reconstruction of a simulated FMC of a point reflector.
 482 Element positions are displayed as circles, while the paths of travel for two pixels at (a, b)
 483 and (c, d) are shown for the same transmission and reception. Part (b) displays a SAFT
 484 reconstruction of the same data, showing the reconstructing triangles for two pixels.

485

486 This principle is repeated over the entire region shown as non-zero in Fig. 16 (a). Because
 487 reflections are not represented in A-scans as infinitely thin peaks, but instead have a width,
 488 the pixels which receive a non-zero contribution need not possess the exact time of travel as
 489 the position of the true reflector. Instead, the time of travel must only be within a range
 490 proximal to the true reflector. This is illustrated in Fig. 17, which shows the exact time of
 491 travel of the point reflector, P, along with the region in which the contribution will be non-
 492 zero, R. This leads to numerous non-zero contributions throughout the image, confined by the
 493 geometrical combinations of pixel position, transmit-receive pairings and reflector position.
 494 There are particular regions where more of the transmit-receive pairs and pixel positions meet
 495 the criteria for non-zero contribution. This leads to regions where the side lobes have a
 496 particularly high intensity, as can be seen in Fig. 16 (a).



497

498 Fig. 17: An A-scan from the point reflector FMC, showing the time of travel for the point
 499 reflector, P, and the range in which pixel contributions will be non-zero, R.

500

501 The 20° SAFT reconstruction of the same FMC data is presented in Fig. 16 (b), displaying a
 502 significantly lower side lobe contribution. Also shown are the reconstructing triangles for two
 503 pixels, with the base determining the receiving elements employed reconstruction sub
 504 aperture. The left pixel, showing a yellow triangle, has included a number of elements in the
 505 reconstructing sub aperture, while the right pixel has none. The result of this is that the left
 506 pixel has a non-zero intensity, while the right has zero intensity.

507 While this example employed a point reflector for simplicity, the principle presented is valid
508 for any physical reflector which elicits a high intensity response in a number of the receiving
509 elements and, therefore, produces a region of high intensity in the reconstructed image. As
510 such, the flat surface shown in Fig. 7 can be thought of as a densely populated line of
511 discrete, strong reflections, each creating side lobes. These have merged to form a thick line
512 above the true surface. This effect is significantly lessened in the SAFT images, as the
513 number of possible transmit-receive pairs is limited by the fact that the number of elements
514 considered in reception is significantly less than that in TFM. This geometrically restricts the
515 regions in which side lobes can be formed, as demonstrated in Fig. 16 (b). Given that a
516 narrow sub aperture has been employed in this study, the likelihood of side lobes is low. If,
517 however, the sub aperture size was increased, the possibility of side lobes would increase.

518 It is pertinent also in this discussion to address the issue of path geometry used in this study
519 and the likelihood of such path scanning to be employed in a real application on a knee joint.
520 In the present study, the fact that the basic geometry of the test samples was known a-priori
521 allowed an immediate construction of a suitable probe scan path. In the final application, the
522 geometry is unknown and such simple path construction is considered unlikely to produce an
523 outcome on a single pass. It is likely that a 2 (or more) stage scan would be required in
524 practice, with an initial coarse scan used to generate a basic representation of the bone
525 surface. This coarse scan would then be used to construct an optimised scan path around the
526 knee joint for the subsequent high resolution scan.

527

528 **4. Summary and Conclusions**

529 The performance of two synthetic aperture methods to accurately image a number of surfaces
530 corresponding to three precision metallic objects with surfaces including curves, flat surfaces
531 and edges has been presented. The Total Focussing Method (TFM) and the Synthetic
532 Aperture Focussing Method (SAFT) were selected for image reconstruction, as these
533 represented extremes of the imaging approach, employing the full synthetic aperture width
534 and a minimal aperture width, respectively.

535 The metallic samples of known geometries (cuboid, sphere and cylinder) were imaged using
536 a 5 MHz, 128 element 1D phased array, which was manipulated around the samples using a
537 precision robotic positioning system, allowing for repeatable ultrasound derived 3D surface

538 point clouds to be created. A global thresholding technique was presented that allowed
539 extraction of the surface profiles and these were compared with the known geometry samples
540 to provide a quantitative measure of error of 3D surface reconstruction.

541 Producing mean errors of 1.3 mm, 2.9 mm and 2.2 mm using SAFT and 3.7 mm, 3.0 mm and
542 3.1 mm using TFM for the cuboidal, spherical and cylindrical samples respectively, SAFT
543 offered significant improvements in accuracy over TFM. This was a result of improved
544 clarity of surface representations, which allowed for more accurate surface profile extraction.
545 The reduction in the width of the synthetic aperture of SAFT allowed for this, eliminating the
546 side lobes associated with TFM. While the use of SAFT imposed a slight reduction in
547 resolution and coverage, it provided mean errors approaching the resolution of CT - the 'gold
548 standard' in preoperative imaging for robotic knee arthroplasty. Therefore, for unknown
549 surface types, a narrow aperture SAFT is the superior imaging method, indicating that it
550 would provide the most accurate depictions of the complex surfaces in the prescribed
551 biomedical application. Additionally, this result has significant implications for dual-media
552 time of flight correction techniques within NDT, in that the employment of a narrow aperture
553 SAFT could allow for automatic identification of the surfaces of parts with more complex
554 shapes than would be possible with TFM.

555 The results presented in this paper indicate that synthetic aperture methods are capable of
556 highly accurate surface imaging, provided a narrow synthetic aperture is employed. While
557 edges and curved surfaces were responsible for errors, the shapes employed were
558 rudimentary and intended to amplify the associated effects. As such, they serve as a worst
559 case scenario and the effects would be expected to be significantly lessened in bone surface
560 imaging due to the increase in shape complexity and decrease in surface specularly.

561 In addition to the change in the surface type, the final application would also include multiple
562 soft tissue interfaces preceding the bony surface under inspection. However, it is predicted
563 this would not cause serious ill-effect in the resulting images, in that the acoustic impedance
564 mismatch between various soft tissues and water is small relative to that between soft tissue
565 and bone. This, however, can only be confirmed by experiment. Therefore, future
566 investigations should concentrate on real bone surfaces with preceding soft tissue layers with
567 an aim to achieve a closer representation of the bony surfaces found in the knee joint.

568

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571

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